

EE 570: Location and Navigation

Navigation Mathematics: Angular and Linear Velocity

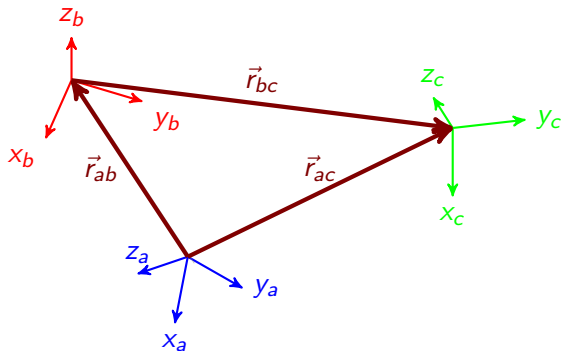
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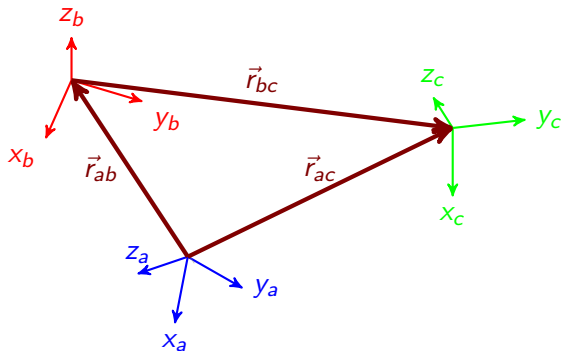
February 10, 2016

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- translation between frames $\{a\}$ and $\{c\}$:

$$\vec{r}_{ac} = \vec{r}_{ab} + \vec{r}_{bc}$$



- translation between frames {a} and {c}:

$$\vec{r}_{ac} = \vec{r}_{ab} + \vec{r}_{bc}$$

- written wrt/frame {a}

$$\begin{aligned} \vec{r}_{ac}^a &= \vec{r}_{ab}^a + \vec{r}_{bc}^a \\ &= \vec{r}_{ab}^a + C_b^a \vec{r}_{bc}^b \end{aligned}$$

- Given relationship for translation between moving (rotating and translating) frames

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what is linear velocity between frames?

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- Why is $\dot{C}_b^a \neq 0$ in general? Re-coordinatization of \vec{r}_{bc}^b is time-dependent.
- \dot{C}_b^a is directly related to angular velocity between frames $\{a\}$ and $\{b\}$.

Given a rotation matrix C , one of its properties is

$$[C_b^a]^T C_b^a = C_b^a [C_b^a]^T = \mathcal{I}$$

Taking the time-derivative of the “right-inverse” property

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$$\qquad \qquad \qquad \underbrace{\hspace{10em}}_{[\Omega_{ab}^a]^T}$$

$$\Rightarrow \Omega_{ab}^a + [\Omega_{ab}^a]^T = 0$$

$\Rightarrow \Omega_{ab}^a$ is skew-symmetric!

Define this skew-symmetric matrix Ω_{ab}^a

$$\Omega_{ab}^a = [\vec{\omega}_{ab}^a \times] = \begin{bmatrix} 0 & -\omega_z & \omega_y \\ \omega_z & 0 & -\omega_x \\ -\omega_y & \omega_x & 0 \end{bmatrix}$$

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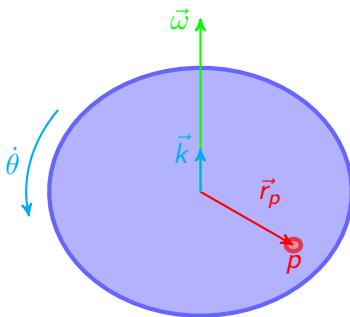
Note $\Omega_{ab}^a = \dot{C}_b^a [C_b^a]^T$

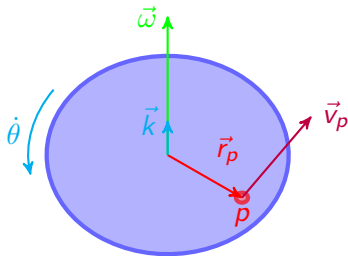
$$\Rightarrow \dot{C}_b^a = \Omega_{ab}^a C_b^a$$

is a means of finding derivative of rotation matrix provided we can further understand Ω_{ab}^a .

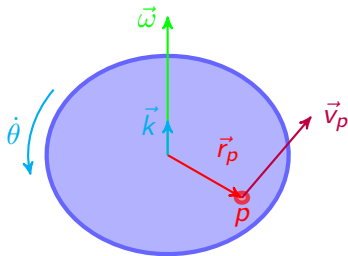
Now for some insight into physical meaning of Ω_{ab}^a .

- Consider a point p on a rigid body rotating with angular velocity $\vec{\omega} = [\omega_x, \omega_y, \omega_z]^T = \dot{\theta} \vec{k} = \dot{\theta} [k_x, k_y, k_z]^T$ with \vec{k} a unit vector.



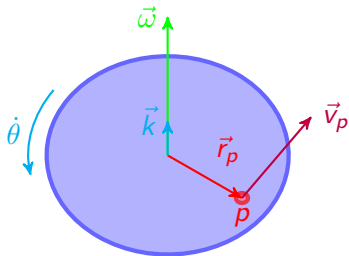


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$$\vec{v}_p = \vec{\omega} \times \vec{r}_p = \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix} \times \begin{bmatrix} r_x \\ r_y \\ r_z \end{bmatrix} = \begin{bmatrix} \omega_y r_z - \omega_z r_y \\ \omega_z r_x - \omega_x r_z \\ \omega_x r_y - \omega_y r_x \end{bmatrix} = \underbrace{\begin{bmatrix} 0 & -\omega_z & \omega_y \\ \omega_z & 0 & -\omega_x \\ -\omega_y & \omega_x & 0 \end{bmatrix}}_{?} \begin{bmatrix} r_x \\ r_y \\ r_z \end{bmatrix}$$

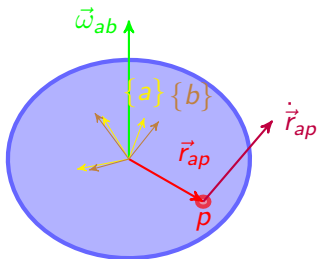


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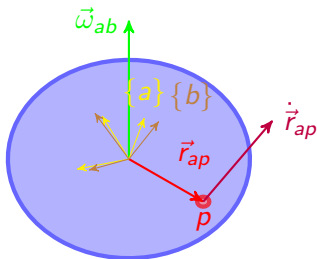
$$\vec{v}_p = \vec{\omega} \times \vec{r}_p = \begin{bmatrix} \omega_x \\ \omega_y \\ \omega_z \end{bmatrix} \times \begin{bmatrix} r_x \\ r_y \\ r_z \end{bmatrix} = \begin{bmatrix} \omega_y r_z - \omega_z r_y \\ \omega_z r_x - \omega_x r_z \\ \omega_x r_y - \omega_y r_x \end{bmatrix} = \underbrace{\begin{bmatrix} 0 & -\omega_z & \omega_y \\ \omega_z & 0 & -\omega_x \\ -\omega_y & \omega_x & 0 \end{bmatrix}}_{\Omega = [\vec{\omega} \times]} \begin{bmatrix} r_x \\ r_y \\ r_z \end{bmatrix}$$

$\Rightarrow \Omega$ represents angular velocity and performs cross product

Now let's add fixed frame $\{a\}$ and rotating frame $\{b\}$ attached to moving body such that there is angular velocity $\vec{\omega}_{ab}$ between them.



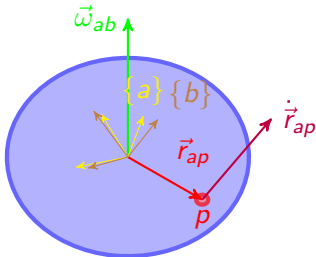
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Start with position

$$\vec{r}_{ap}^a = \underbrace{\vec{r}_{ab}^a}_0 + C_b^a \vec{r}_{bp}^b$$

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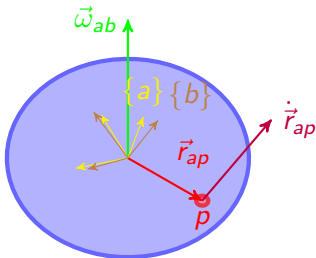
and take derivative wrt time

$$\begin{aligned} \dot{\vec{r}}_{ap}^a &= \underbrace{\dot{C}_b^a}_{\Omega_{ab}^a C_b^a} \vec{r}_{bp}^b + \underbrace{C_b^a \dot{\vec{r}}_{bp}^b}_0 \\ &= \Omega_{ab}^a C_b^a \vec{r}_{bp}^b \\ &= \Omega_{ab}^a \vec{r}_{ap}^a = [\vec{\omega}_{ab}^a \times] \vec{r}_{ap}^a \end{aligned}$$

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from which it is observed (compare to $\vec{v}_p = \vec{\omega} \times \vec{r}_p$) that Ω_{ab}^a represents cross product with angular velocity $\vec{\omega}_{ab}^a$.

- Another approach to developing derivative of rotation matrix and angular velocity is based upon angle-axis representation of orientation and rotation matrix as exponential.
- This approach is included in notes.

$$\begin{aligned}
 C\Omega C^T \vec{b} &= C \left[\vec{\omega} \times (C^T \vec{b}) \right] \\
 &= C \vec{\omega} \times (CC^T \vec{b}) \\
 &= C \vec{\omega} \times \vec{b} \\
 &= [C\vec{\omega} \times] \vec{b}
 \end{aligned}$$

Therefore (from above),

$$C\Omega C^T = C[\vec{\omega} \times]C^T = [C\vec{\omega} \times]$$

and (via distributive property)

$$C[\vec{\omega} \times] = [C\vec{\omega} \times]C$$

$$\begin{aligned}\dot{C}_b^a &= \Omega_{ab}^a C_b^a \\ &= [\vec{\omega}_{ab}^a \times] C_b^a \\ &= [C_b^a \vec{\omega}_{ab}^b \times] C_b^a \\ &= C_b^a [\vec{\omega}_{ab}^b \times] \\ &= C_b^a \Omega_{ab}^b\end{aligned}$$

$$\Rightarrow \dot{C}_b^a = \Omega_{ab}^a C_b^a = C_b^a \Omega_{ab}^b$$

Angular velocity can be

- described as a vector
 - the angular velocity of the b -frame wrt the a -frame resolved in the c -frame, $\vec{\omega}_{ab}^c$
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- described as a skew-symmetric matrix $\Omega_{ab}^c = [\vec{\omega}_{ab}^c \times]$
 - the skew-symmetric matrix is equivalent to the vector cross product when pre-multiplying another vector
- related to the derivative of the rotation matrix

$$\dot{C}_b^a = \Omega_{ab}^a C_b^a = C_b^a \Omega_{ab}^b$$

$$\dot{C}_b^a = -\Omega_{ba}^a C_b^a = -C_b^a \Omega_{ba}^b$$

Consider the derivative of the composition of rotations $C_2^0 = C_1^0 C_2^1$.

$$\frac{d}{dt} C_2^0 = \frac{d}{dt} C_1^0 C_2^1$$

$$\dot{C}_2^0 = \dot{C}_1^0 C_2^1 + C_1^0 \dot{C}_2^1$$

$$\Omega_{02}^0 C_2^0 = \Omega_{01}^0 C_1^0 C_2^1 + C_1^0 C_2^1 \Omega_{12}^2$$

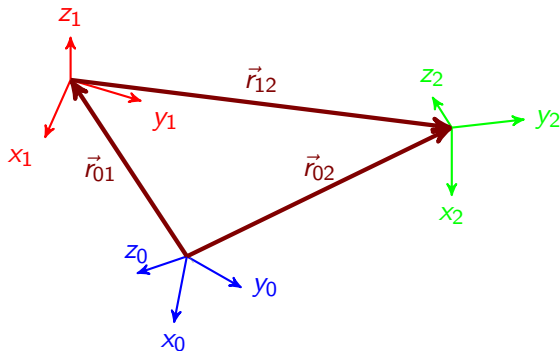
$$\Omega_{02}^0 = \Omega_{01}^0 C_2^0 [C_2^0]^T + C_2^0 \Omega_{12}^2 [C_2^0]^T$$

$$[\vec{\omega}_{02}^0 \times] = [\vec{\omega}_{01}^0 \times] + [C_2^0 \vec{\omega}_{12}^2 \times]$$

$$\Rightarrow \vec{\omega}_{02}^0 = \vec{\omega}_{01}^0 + \vec{\omega}_{12}^0$$

\Rightarrow angular velocities (as vectors) add so long as resolved common coordinate system

We can get back to where we started ... motion (translation and rotation) between frames and their derivatives.



Translation (position) between frames $\{0\}$ and $\{1\}$:

$$\begin{aligned}\vec{r}_{02}^0 &= \vec{r}_{01}^0 + \vec{r}_{12}^0 \\ &= \vec{r}_{01}^0 + C_1^0 \vec{r}_{12}^1\end{aligned}$$

Linear velocity:

$$\begin{aligned}
 \dot{\vec{r}}_{02}^0(t) &= \frac{d}{dt} (\vec{r}_{01}^0 + C_1^0 \vec{r}_{12}^1) \\
 &= \dot{\vec{r}}_{01}^0 + \dot{C}_1^0 \vec{r}_{12}^1 + C_1^0 \dot{\vec{r}}_{12}^1 \\
 &= \dot{\vec{r}}_{01}^0 + [\vec{\omega}_{01}^0 \times] C_1^0 \vec{r}_{12}^1 + C_1^0 \dot{\vec{r}}_{12}^1 \\
 &= \dot{\vec{r}}_{01}^0 + \vec{\omega}_{01}^0 \times (C_1^0 \vec{r}_{12}^1) + C_1^0 \dot{\vec{r}}_{12}^1
 \end{aligned}$$

Linear acceleration:

$$\begin{aligned} \ddot{\mathbf{r}}_{02}^0 &= \frac{d}{dt} \left(\dot{\mathbf{r}}_{01}^0 + \bar{\omega}_{01}^0 \times \left(C_1^0 \mathbf{r}_{12}^1 \right) + C_1^0 \dot{\mathbf{r}}_{12}^1 \right) \\ &= \ddot{\mathbf{r}}_{01}^0 + \dot{\bar{\omega}}_{01}^0 \times \left(C_1^0 \mathbf{r}_{12}^1 \right) + \bar{\omega}_{01}^0 \times \left(\dot{C}_1^0 \mathbf{r}_{12}^1 \right) + \bar{\omega}_{01}^0 \times \left(C_1^0 \dot{\mathbf{r}}_{12}^1 \right) + \dot{C}_1^0 \dot{\mathbf{r}}_{12}^1 + C_1^0 \ddot{\mathbf{r}}_{12}^1 \end{aligned}$$

$$= \ddot{\mathbf{r}}_{01}^0 + \dot{\bar{\omega}}_{01}^0 \times \mathbf{r}_{12}^0(t) + \bar{\omega}_{01}^0 \times \left(\bar{\omega}_{01}^0 \times \mathbf{r}_{12}^0(t) \right) + 2\bar{\omega}_{01}^0 \times \left(C_1^0 \dot{\mathbf{r}}_{12}^1 \right) + C_1^0 \ddot{\mathbf{r}}_{12}^1$$

accel of {1}'s origin
from {0} in {0}

Centripetal accel ($\omega^2 r$)

Coriolis accel ($2\omega \times v$)

Transverse accel

accel of {2}'s origin
from {1} in {0}

